## Batalin-Vilkovisky formalism beyond perturbation theory via derived geometry

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Talk at
Texas Tech Topology and Geometry Seminar Texas Tech University \& Wichita State University, Kansas
19/03/2024

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### 1.1 BV-theory

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This is hard, and problematic if there is a gauge symmetry.

## Idea of BV-theory

Look at its derived critical locus $\mathbb{R C r i t}(S)$, a derived enhancement of $\operatorname{Crit}(S)$.

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## Idea of BV-theory

Look at its derived critical locus $\mathbb{R} \operatorname{Crit}(S)$, a derived enhancement of $\operatorname{Crit}(S)$.

Informally speaking,

$$
\mathbb{R} \operatorname{Crit}(S)=\left\{\phi \in \underset{\text { space }}{\text { configuration }}, \phi^{+} \in \text { antifields } \mid \delta S(\phi) \xrightarrow{\phi^{+}} 0\right\}
$$

E.o.m. not imposed on the nose, but up to something, a 1-simplex.

### 1.1 BV-theory

Main approaches to make classical (and quantum) BV-theory precise in the literature:
(1) NQP-manifolds/ $L_{\infty}$-algebroid approach. [Jurčo, Raspollini, Sämann, Wolf, ...] Algebra of classical observables is given by a Poisson dg-Lie algebra of functions on an NQP-manifold, i.e. a differential-graded manifold (dg-manifold) equipped with a ( -1 )-shifted symplectic form. (Equivalently, a symplectic $L_{\infty}$-algebroid.)

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(2) (Perturbative) Algebraic Quantum Field Theory. [Schenkel, Benini, Rejzner, ...] Algebra of observables is given by a net of differential-graded Poisson *-algebras, on spacetime.
(3) Factorisation Algebras approach. [Costello, Gwilliam, Williams, ...]

Algebra of classical observables is given by the differential-graded $\mathbb{P}_{0}$-algebra of functions on a ( -1 )-shifted symplectic formal moduli problem, which is sheaved on spacetime.

Approaches (1) \& (3) very close. Approaches (2) \& (3) related by [Schenkel, Benini, ...].

### 1.2 Familiar recipe for BV-theory

- Ingredients:
(1) an $L_{\infty}$-algebra $\mathfrak{L}$ (BRST-algebra),
(2) an element $S \in \mathrm{CE}(\mathfrak{L})$ (action functional),
where:

$$
\mathrm{CE}(\mathfrak{L})=\left(\operatorname{Sym} \mathfrak{L}^{\vee}[-1], \mathrm{d}_{\mathrm{CE}(\mathfrak{L})}\right)
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- The ( -1 )-shifted cotangent bundle

$$
T^{\vee}[-1] \mathfrak{L}[1]=\left(\mathfrak{L} \oplus \mathfrak{L}^{\vee}[-3]\right)[1] .
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comes with natural $(-1)$-shifted Poisson bracket $\{-,-\}$.

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- Construct $S_{\mathrm{BV}}:=S+S_{\mathrm{BRST}}(B V$-action $)$ where $S_{\mathrm{BRST}}:=\widehat{\mathrm{d}_{\mathrm{CE}(\mathfrak{L})}}$.
- Since $\left\{S_{\mathrm{BV}}, S_{\mathrm{BV}}\right\}=0$, we have a new $L_{\infty}$-algebra $\mathfrak{C r i t}(S)$ given by

$$
\operatorname{CE}(\mathfrak{C r i t}(S)):=\left(\operatorname{Sym}\left(\mathfrak{L}^{\vee}[-1] \oplus \mathfrak{L}[2]\right), \quad Q_{\mathrm{BV}}=\left\{S_{\mathrm{BV}},-\right\}\right)
$$

This is the BV-complex.

## Example: Yang-Mills theory

- BRST algebra

$$
\begin{array}{rl}
\mathfrak{L}[1]=\left(\Omega^{0}(M, \mathfrak{g}) \xrightarrow{\mathrm{d}} \Omega^{1}(M, \mathfrak{g})\right) \\
\operatorname{deg}=\quad-1 & 0 \\
\ell_{1}(c) & =\mathrm{d} c \\
\ell_{2}\left(c_{1}, c_{2}\right) & =\left[c_{1}, c_{2}\right]_{\mathfrak{g}}, \\
\ell_{2}(c, A) & =[c, A]_{\mathfrak{g}},
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- BV-BRST algebra

$$
\begin{gathered}
\operatorname{crit}(S)[1]=\left(\Omega^{0}(M, \mathfrak{g}) \xrightarrow{\mathrm{d}} \Omega^{1}(M, \mathfrak{g}) \xrightarrow{\mathrm{d} \star \mathrm{~d}} \Omega^{d-1}(M, \mathfrak{g}) \xrightarrow{\mathrm{d}} \Omega^{d}(M, \mathfrak{g})\right) \\
\operatorname{deg}=\quad{ }_{-1} \\
S_{\mathrm{BV}}\left(\mathrm{c}, \mathrm{~A}, \mathrm{~A}^{+}, \mathrm{c}^{+}\right)=\int_{M}(\underbrace{\left(\frac{1}{2}\left\langle F_{\mathrm{A}}, \star F_{\mathrm{A}}\right\rangle_{\mathfrak{g}}\right.}_{\mathrm{S}}-\underbrace{\left\langle\mathrm{A}^{+}, \nabla_{\mathrm{A} C}\right\rangle_{\mathfrak{g}}+\frac{1}{2}\left\langle\mathrm{c}^{+},[\mathrm{c}, \mathrm{c}]_{\mathfrak{g}}\right\rangle_{\mathfrak{g}}}_{S_{\mathrm{BRST}}}) .
\end{gathered}
$$

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- BV-BRST algebra

$$
\begin{aligned}
& \mathfrak{C r i t}(S)[1]=\left(\Omega^{0}(M, \mathfrak{g}) \xrightarrow{\mathrm{d}} \Omega^{1}(M, \mathfrak{g}) \xrightarrow{\mathrm{d} \star \mathrm{~d}} \Omega^{d-1}(M, \mathfrak{g}) \xrightarrow{\mathrm{d}} \Omega^{d}(M, \mathfrak{g})\right) \\
& \begin{array}{lllll}
\operatorname{deg}= & -1 & 0 & 1 & 2
\end{array} \\
& \ell_{1}(c)=\mathrm{d} c, \\
& \ell_{1}(A)=\mathrm{d} \star \mathrm{~d} A, \quad \quad \ell_{1}\left(A^{+}\right)=\mathrm{d} A^{+}, \\
& \ell_{2}\left(c_{1}, c_{2}\right)=\left[c_{1}, c_{2}\right]_{\mathfrak{g}}, \quad \ell_{2}\left(c, c^{+}\right)=\left[c, c^{+}\right]_{\mathfrak{g}}, \\
& \ell_{2}(c, A)=[c, A]_{\mathfrak{g}}, \quad \ell_{2}\left(c, A^{+}\right)=\left[c, A^{+}\right]_{\mathfrak{g}}, \\
& \ell_{2}\left(A, A^{+}\right)=\left[A \hat{,} A^{+}\right]_{\mathfrak{g}}, \\
& \ell_{2}\left(A_{1}, A_{2}\right)=\mathrm{d} \star\left[A_{1} \hat{,} A_{2}\right]_{\mathfrak{g}}+\left[A_{1} \hat{,} \star \mathrm{~d} A_{2}\right]_{\mathfrak{g}}+\left[A_{2} \hat{,} \star \mathrm{~d} A_{1}\right]_{\mathfrak{g}}, \\
& \ell_{3}\left(A_{1}, A_{2}, A_{3}\right)=\left[A_{1} \hat{,} \star\left[A_{2} \hat{,} A_{3}\right]_{\mathfrak{g}}\right]_{\mathfrak{g}}+\left[A_{2} \hat{,} \star\left[A_{3} \hat{,} A_{1}\right]_{\mathfrak{g}}\right]_{\mathfrak{g}}+\left[A_{3} \hat{,} \star\left[A_{1} \hat{,} A_{2}\right]_{\mathfrak{g}}\right]_{\mathfrak{g}},
\end{aligned}
$$

### 1.3 BV-theory as deformation theory

- Artinian dg-algebra: finite-dimensional, non-positively graded, dg-commutative algebra $\mathcal{R}$ s.t. it has a unique maximal differential ideal $\mathfrak{m}_{\mathcal{R}}$ which is nilpotent and $\mathcal{R} / \mathfrak{m}_{\mathcal{R}} \cong \mathbb{R}$.
- Formal Moduli Problem: (algebraic) stack on Artinian dg-algebras, i.e.

$$
F: \operatorname{dgArt}{ }^{\leq 0} \longrightarrow \text { sSet. }
$$

- Any formal moduli problem is equivalent to $F \simeq \operatorname{MC}(\mathfrak{g})$, for some $L_{\infty}$-algebra $\mathfrak{g}$, where

$$
\mathrm{MC}(\mathfrak{g}): \mathcal{R} \longmapsto \mathrm{MC}\left(\mathfrak{g} \otimes \mathfrak{m}_{\mathcal{R}}\right) .
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In BV-theory:
MC(Crit(S)) is the derived critical locus of the action functional $S$ on $\mathbf{M C}(\mathfrak{L})$

## Example: Yang-Mills theory

How does the story go for Yang-Mills theory?
$\operatorname{MC}(\mathfrak{C r i t}(S)): \mathcal{R} \longmapsto \operatorname{MC}\left(\mathfrak{C r i t}(S) \otimes \mathfrak{m}_{\mathcal{R}}\right)$

## Example: Yang-Mills theory

## How does the story go for Yang-Mills theory?

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\operatorname{MC}(\mathfrak{C r i t}(S)): \mathcal{R} \longmapsto \operatorname{MC}\left(\mathfrak{C r i t}(S) \otimes \mathfrak{m}_{\mathcal{R}}\right)
$$

$$
\begin{aligned}
& \operatorname{MC}\left(\mathfrak{C r i t}(S) \otimes \mathfrak{m}_{\mathcal{R}}\right)_{0}=\left\{\begin{array}{l|r}
A & \in \Omega^{1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R}, 0} \\
A^{+} \in \Omega^{d-1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-1} & \nabla_{A} \star F_{A}=\mathrm{d}_{\mathcal{R}} A^{+} \\
c^{+} \in \Omega^{d}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-2} & \nabla_{A} A^{+}=\mathrm{d}_{\mathcal{R}} c^{+}
\end{array}\right\}, \\
& \operatorname{MC}\left(\mathfrak{C r i t}(S) \otimes \mathfrak{m}_{\mathcal{R}}\right)_{1}=\left\{\begin{array}{llr}
c_{1} \mathrm{~d} t \in \Omega^{0}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R}, 0} \otimes \Omega^{1}([0,1]) & \nabla_{A_{0}} \star F_{A_{0}}=\mathrm{d}_{\mathcal{R}} A_{0}^{+} \\
A_{0} & \in \Omega^{1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R}, 0} \otimes \Omega^{0}([0,1]) & \nabla_{A_{0}} A_{0}^{+}=\mathrm{d}_{\mathcal{R}} c_{0}^{+} \\
A^{1} \mathrm{~d} t \in \Omega^{1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-1} \otimes \Omega^{1}([0,1]) & \frac{\mathrm{d}}{\mathrm{~d} t} A_{0}+\nabla_{A_{0}} c_{1}=\mathrm{d}_{\mathcal{R}} A_{1} \\
A_{0}^{+} & \in \Omega^{d-1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-1} \otimes \Omega^{0}([0,1]) & \frac{\mathrm{d}}{\mathrm{~d} t} A_{0}^{+}+\nabla_{A_{0}} \star F_{A_{1}}+ \\
A_{1}^{+} \mathrm{d} t \in \Omega^{d-1}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-2} \otimes \Omega^{1}([0,1]) & +\left[c_{1}, A_{0}^{+}=\mathrm{d}_{\mathcal{R}} A_{1}^{+}\right. \\
c_{0}^{+} & \in \Omega^{d}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-2} \otimes \Omega^{0}([0,1]) & \frac{\mathrm{d}}{\mathrm{~d} t} c_{0}^{+}+\nabla_{A_{0}} A_{1}^{+}+ \\
c_{1}^{+} \mathrm{d} t \in \Omega^{d}(M, \mathfrak{g}) \otimes \mathfrak{m}_{\mathcal{R},-3} \otimes \Omega^{1}([0,1]) & +\left[c_{1}, c_{0}^{+}\right]=\mathrm{d}_{\mathcal{R}} c_{1}^{+}
\end{array}\right\},
\end{aligned}
$$

$$
\left(c_{1} \mathrm{~d} t, A_{0}+A_{1} \mathrm{~d} t, A_{0}^{+}+A_{1}^{+} \mathrm{d} t, c_{0}^{+}+c_{1}^{+} \mathrm{d} t\right)
$$

$$
\left(A, A^{+}, c^{+}\right) \quad\left(A^{\prime}, A^{+\prime}, c^{+\prime}\right)
$$

### 1.4 Global BV-theory?

Formal Moduli Problem: (algebraic) derived stack on Artinian dg-algebras, i.e.

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F: \operatorname{dgArt}^{\leq 0} \longrightarrow \text { sSet }
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Artinian dg-algebras $\simeq$ algebras of function on "derived thickened points".

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A ( -1 )-symplectic Formal Moduli Problem can be seen as the formal completion of a fully-fledged ( -1 )-symplectic derived stack at some given point.


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We have the following picture:

> Formal Moduli Problem $\longleftrightarrow$ Perturbative physics
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In fact, in terms of configuration spaces:

## Example (Stack of $G$-bundles with connection)

$$
\underbrace{\left[\Omega^{1}(M, \mathfrak{g}) / \mathcal{C}^{\infty}(M, \mathfrak{g})\right]}_{L_{\infty} \text {-algebroid }} \neq \underbrace{\operatorname{Bun}_{G}^{\nabla}(M)}_{\text {stack of } G \text {-bundles }}:=\left[M, \mathbf{B} G_{\text {conn }}\right]
$$

Physics includes (higher) gauge theories

- Quantisation requires BV-theory, i.e. derived geometry
- Finite (higher) gauge transformations and global properties require stacks, i.e. higher geometry (e.g. Aharonov-Bohm phase and magnetic charge for electromagnetic field)


### 1.5 Smooth stacks

smooth manifolds

### 1.6 Global BRST formalism

- An ordinary geometric space can be encoded by its functor of points, which is an ordinary sheaf.



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- An ordinary geometric space can be encoded by its functor of points, which is an ordinary sheaf.

- A higher geometric space can be defined as a stack, which is a functor

$$
X: \mathrm{Mfd}^{\mathrm{op}} \longrightarrow \mathrm{sSet}
$$

satisfying a higher sheaf condition, i.e. it is an element of

$$
\text { SmoothStack }:=\left[\mathrm{Mfd}^{\mathrm{op}}, \mathrm{sSet}\right]_{\text {proj, loc }}^{\circ}
$$

$\Longrightarrow$ geometry encompassing gauge principle from physics.

### 1.5 Global BRST formalism

Now, let us go back to smooth stacks.
Moduli stack of principal $G$-bundles:

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\mathbf{B} G=[* / G]
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$$

Let $M$ be a smooth manifold and $\coprod_{\alpha \in I} V_{\alpha} \rightarrow M$ be a good open cover of it.
A map $M \xrightarrow{g_{\alpha \beta}} \mathbf{B} G$ is given by a $G$-bundle $\left\{g_{\alpha \beta} \in \mathcal{C}^{\infty}\left(V_{\alpha} \cap V_{\beta}, G\right) \mid g_{\alpha \beta} \cdot g_{\beta \gamma}=g_{\alpha \gamma}\right\}$.

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Smooth stack of principal $G$-bundles on the manifold $M$ :

$$
\operatorname{Bun}_{G}(M):=[M, \mathbf{B} G] .
$$

### 1.6 Global BRST formalism

Smooth stack $\operatorname{Bun}_{G}^{\nabla}(M)$ of principal $G$-bundles with connection:

$$
\operatorname{Hom}\left(U, \operatorname{Bun}_{G}^{\nabla}(M)\right) \simeq \operatorname{cosk}_{2}\left(Z_{2} \underset{\substack{-\left(c_{\alpha}^{\prime}, c_{\alpha \beta}^{\prime}, \begin{array}{c}
\left.g_{\alpha \beta}, A_{\alpha} \\
g_{\alpha \beta}^{\prime \prime}, A_{\alpha}^{\prime \prime}\right) \longrightarrow \\
g_{\alpha \beta}^{\prime \prime}, A_{\alpha}^{\prime \prime}
\end{array}\right)}}{\underset{\left(g_{\alpha \beta}^{\prime}, A_{\alpha}^{\prime}\right)}{\binom{g_{\alpha \beta}, A_{\alpha}}{g_{\alpha \beta}^{\prime}, A_{\alpha}^{\prime}}}} Z_{1} \xrightarrow{\left(g_{\alpha \beta}^{\prime}, A_{\alpha}\right)} Z_{0}\right)
$$

where:

$$
\begin{aligned}
& Z_{0}=\left\{\begin{array}{l|l}
g_{\alpha \beta} \in \mathcal{C}^{\infty}\left(V_{\alpha} \cap V_{\beta} \times U, G\right) & g_{\alpha \beta} \cdot g_{\beta \gamma} \cdot g_{\gamma \alpha}=1 \\
A_{\alpha} \in \Omega_{\mathrm{ver}}^{1}\left(V_{\alpha} \times U, \mathfrak{g}\right) & A_{\alpha}=g_{\beta \alpha}^{-1}\left(A_{\beta}+\mathrm{d}\right) g_{\beta \alpha}
\end{array}\right\}, \\
& Z_{1}=\left\{\begin{array}{l|l|l} 
& g_{\alpha \beta} \cdot g_{\beta \gamma} \cdot g_{\gamma \alpha}=1 \\
c_{\alpha} & \in \mathcal{C}^{\infty}\left(V_{\alpha} \times U, G\right) & A_{\alpha}=g_{\beta \alpha}^{-1}\left(A_{\beta}+\mathrm{d}\right) g_{\beta \alpha} \\
g_{\alpha \beta}, g_{\alpha \beta}^{\prime} \in \mathcal{C}^{\infty}\left(V_{\alpha} \cap V_{\beta} \times U, G\right) & g_{\alpha \beta}^{\prime} \cdot g_{\beta \gamma}^{\prime} \cdot g_{\gamma \alpha}^{\prime}=1 \\
A_{\alpha}, A_{\alpha}^{\prime} & \in \Omega_{\mathrm{ver}}^{1}\left(V_{\alpha} \times U, \mathfrak{g}\right) & A_{\alpha}^{\prime}=g_{\beta \alpha}^{\prime-1}\left(A_{\beta}^{\prime}+\mathrm{d}\right) g_{\beta \alpha}^{\prime} \\
& g_{\alpha \beta}^{\prime}=c_{\beta}^{-1} g_{\alpha \beta} c_{\alpha} \\
A_{\alpha}^{\prime}=c_{\alpha}^{-1}\left(A_{\alpha}+\mathrm{d}\right) c_{\alpha}
\end{array}\right\},
\end{aligned}
$$

$Z_{2}=$ \{compositions of gauge transformations $\}$,

### 2.1 Family tree of smooth stacks

smooth manifolds

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### 2.3 Derived smooth manifolds

- Observation: given manifolds $M, N \hookrightarrow B$, the intersection $M \cap N:=M \times_{B} N$ is not generally well-defined in Mfd.
- Solution: derived smooth manifolds [Spivak, Joyce, Carchedi, Steffens, ...].


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- Solution: derived smooth manifolds [Spivak, Joyce, Carchedi, Steffens, ...].

Use the natural embedding:

$$
i: \mathbf{N}(\mathrm{Mfd}) \longrightarrow \mathbf{d M f d}
$$

The derived intersection always exists in the ( $\infty, 1$ )-category dMfd:



For example, in $B=\mathbb{R}^{3}$,

$$
\begin{aligned}
& M=\left\{\left(x, y, x^{2} y^{2}\right) \in \mathbb{R}^{3} \mid(x, y) \in \mathbb{R}^{2}\right\} \\
& \text { and } N=x, y \text {-plane. }
\end{aligned}
$$

## $2.2 \mathcal{C}^{\infty}$-algebras

Let CartSp be the category of Cartesian spaces $\left\{\mathbb{R}^{n}\right\}_{n \in \mathbb{N}}$ and smooth maps between them.
This is a Lawvere theory, as any object is such that $\mathbb{R}^{n} \cong \mathbb{R} \times \cdots \times \mathbb{R}$.

## Definition

A $\mathcal{C}^{\infty}$-algebra is a product-preserving functor

$$
A: \text { CartSp } \longrightarrow \text { Set. }
$$

## Example

Let $M \in$ Mfd be a smooth manifold.

$$
\mathcal{C}^{\infty}(M): \mathbb{R}^{n} \longmapsto \mathcal{C}^{\infty}\left(M, \mathbb{R}^{n}\right)
$$

There is a natural embedding:

$$
\mathrm{Mfd} \hookrightarrow \mathrm{C}^{\infty} \mathrm{Alg}^{\mathrm{op}}
$$

### 2.3 Derived smooth manifolds

Homotopy $\mathcal{C}^{\infty}$-algebras: simplicial $\mathcal{C}^{\infty}$-algebras with projective model structure, i.e.

$$
\mathbf{h C ^ { \infty }} \mathbf{A l g}:=\mathbf{N}_{h c}\left(\left[\Delta^{\mathrm{op}}, \mathrm{C}^{\infty} \mathrm{Alg}\right]_{\mathrm{proj}}^{\circ}\right),
$$

where $\Delta$ is the simplex category.

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$$

where $\Delta$ is the simplex category.
The following will be our effective definition of formal derived manifolds.

## Theorem [Carchedi, Steffens 2019]

There is a canonical equivalence of $(\infty, 1)$-categories

$$
\mathbf{d M f d} \simeq \mathbf{h C}^{\infty} \mathbf{A l g}_{\mathrm{fp}}^{\mathrm{op}}
$$

between the $(\infty, 1)$-category of derived manifolds, and the opposite of the $(\infty, 1)$-category of homotopically finitely presented homotopy $\mathcal{C}^{\infty}$-algebras.

At an intuitive level, $U \in \mathbf{d M f d}$ is a geometric object whose algebra of smooth function is a homotopically finitely presented homotopy $\mathcal{C}^{\infty}$-algebra modelled as

$$
\mathcal{O}(U)=\left(\begin{array}{ll} 
& \rightleftarrows \\
\cdots & \mathcal{Z}(U)_{3} \\
& \rightrightarrows \mathcal{O}(U)_{2} \\
& \left.\longrightarrow \mathcal{O}(U)_{1} \longrightarrow \mathcal{O}(U)_{0}\right)
\end{array}\right.
$$

where each $\mathcal{O}(U)_{i}$ is an ordinary $\mathcal{C}^{\infty}$-algebra.

### 2.4 Formal derived smooth manifolds

Derived smooth manifold do not include objects like

$$
\operatorname{Spec}\left(\frac{\mathcal{C}^{\infty}(\mathbb{R})}{\left(x^{2}\right)}\right)
$$

## Definition

A homotopy $\mathcal{C}^{\infty}$-algebra $A$ is finitely generated if $\pi_{0} A$ is finitely generated as an ordinary $\mathcal{C}^{\infty}$-algebra, i.e. such that $\pi_{0} A \cong \mathcal{C}^{\infty}\left(\mathbb{R}^{n}\right) / \mathcal{I}$.

Let $\mathbf{s} \mathbf{C}^{\infty} \mathbf{A l g}_{\mathrm{fg}}$ be the $(\infty, 1)$-category of finitely generated homotopy $\mathcal{C}^{\infty}$-algebras.

## Definition

We define the $(\infty, 1)$-category of formal derived smooth manifolds by

$$
\mathbf{d F M f d}:=\mathbf{s C}^{\infty} \mathbf{A l g}_{\mathrm{fg}}^{\mathrm{op}}
$$

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### 3.1 Formal derived smooth stacks

- We can define étale maps of formal derived smooth manifolds so that they truncate to ordinary étale maps (they generalise local diffeomorphisms of ordinary manifolds).
- By using étale maps, we can make dFMfd into a $(\infty, 1)$-site.
- By [Toen, Vezzosi 2006], we can define formal derived smooth stacks by

$$
\mathbf{d F S m o o t h S t a c k}:=\left[\mathrm{dFMfd}^{\mathrm{op}}, \mathrm{sSet}\right]_{\text {proj,loc }}^{\circ}
$$

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$$



One has a natural (coreflective) embedding

$$
\text { dFSmoothStack } \underset{t_{0}}{\leftrightarrows} \text { SmoothStack. }
$$

[See Carchedi's and Steffens' current foundational work on derived differential geometry.]

### 3.2 Derived differential geometry

dFSmoothStack comes with differential structure, as defined in [Schreiber 2013]

## de Rham stack

Given a formal derived smooth stack $X$, define

$$
\mathfrak{I}(X): R \longmapsto X\left(R^{\text {red }}\right) \quad \text { with } \quad R^{\text {red }}:=\pi_{0} R / \mathfrak{m}_{\pi_{0} R} .
$$

There is a natural map

$$
\mathfrak{i}_{X}: X \longrightarrow \Im(X)
$$

Similarly to [Khavkine, Schreiber], the differential structure can be used to deal with infinitesimal geometry.

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\mathfrak{i}_{X}: X \longrightarrow \mathfrak{I}(X)
$$

Similarly to [Khavkine, Schreiber], the differential structure can be used to deal with infinitesimal geometry.

A derived infinitesimal disks at $x \in X$ is defined by the pullback

where $\Im(X)$ is the de Rham stack of $X$.
3.3 Formal moduli problems as infinitesimal cohesion

Let FMP be the $(\infty, 1)$-category of Formal Moduli Problems, which can be seen as formal derived stacks on derived infinitesimal disks.


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### 4.2 Derived critical locus of Yang-Mills action

- Yang-Mills action functional as morphism of smooth stacks

$$
\begin{aligned}
S: \operatorname{Bun}_{G}^{\nabla}(M) & \longrightarrow \operatorname{Dens}_{M} \\
\left(g_{\alpha \beta}, A_{\alpha}\right) & \longmapsto \frac{1}{2}\left\langle F_{A} \hat{,} \star F_{A}\right\rangle_{\mathfrak{g}}
\end{aligned}
$$

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\end{aligned}
$$

- Yang-Mills e.o.m as morphism of formal derived smooth stacks

$$
\begin{aligned}
\delta S: \operatorname{Bun}_{G}^{\nabla}(M) & \longrightarrow T_{\text {res }}^{\vee} \operatorname{Bun}_{G}^{\nabla}(M) \\
\left(g_{\alpha \beta}, A_{\alpha}\right) & \longmapsto\left(g_{\alpha \beta}, A_{\alpha}, \nabla_{A_{\alpha} \star} F_{A_{\alpha}}, 0\right)
\end{aligned}
$$

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\end{aligned}
$$

- Construct the derived critical locus of Yang-Mills action



### 4.3 Non-perturbative classical BV-BRST theory

$\mathbb{R} \operatorname{Hom}\left(U, \mathbb{R} \operatorname{Crit}(S)(M)^{\text {pre }}\right) \simeq$


- 0-simplices:
- $g_{\alpha \beta}$ transition functions,
- $A_{\alpha}$ connection,
- $A_{\alpha}^{+}$equations of motion,
- $c_{\alpha}^{+}$Noether identities,
- 1-simplices:
- $c_{\alpha}$ gauge transformations,
- $g_{1, \alpha \beta}$ homotopies of transition functions,
- $A_{1, \alpha}$ homotopies of connections,
- $A_{1, \alpha}^{+}$homotopies of equations of motions,
- $c_{1, \alpha}^{+}$homotopies of Noether identities,
- ( $n \geq 2$ )-simplices: compositions of gauge transformations and homotopies of homotopies.


### 4.4 Recovering usual BV-BRST theory

- Use the de Rham stack to obtain the derived infinitesimal disk of $\mathbb{R} \operatorname{Crit}(S)$ at fixed solution $\left(P, \nabla_{A}\right) \in \mathbb{R} \operatorname{Crit}(S)$ of the e.o.m

- Obtain $L_{\infty}$-algebra with underlying complex

$$
\begin{aligned}
\overrightarrow{\mathfrak{C}^{2} \mathfrak{r i t}}(S)_{\left(P, \nabla_{A}\right)}[1] & =\left(\Omega^{0}\left(M, \mathfrak{g}_{P}\right) \xrightarrow{\nabla_{A}} \Omega^{1}\left(M, \mathfrak{g}_{P}\right) \xrightarrow{\nabla_{A^{\star} \nabla_{A}}^{\longrightarrow}} \Omega^{d-1}\left(M, \mathfrak{g}_{P}\right) \xrightarrow{\nabla_{A}} \Omega^{d}\left(M, \mathfrak{g}_{P}\right)\right) \\
\operatorname{deg} & =\quad-1
\end{aligned}
$$

with expected bracket structure.

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## Outlook: the puzzle of quantisation

n-plectic geometry (or higher symplectic geometry) [Rogers, Baez, Saemann, Szabo, Bunk, Fiorenza, Schreiber, Sati, ...] naturally fits in the following picture:


## Example (Closed string)

[Waldorf 2009]: transgression of a bundle gerbe on a smooth manifold $M$ to a principal $U(1)$-bundle on the loop space $\mathcal{L} M=\left[S^{1}, M\right]$.

## Outlook: the puzzle of quantisation

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- [Ševera 2000]: Courant 2-algebroid and Vinogradov n-algebroid are higher generalisations of the Poisson 1-algebroid (as symplectic $L_{\infty}$-algebroids).
- [Rogers 2011], [Sämann, Ritter 2015]: relation between the $L_{\infty}$-algebras of observables on $n$-plectic manifolds and Vinogradov $n$-algebroids.


## Outlook: the puzzle of quantisation

Topic of current research: derived n-plectic geometry.


Framework to make contact with BV-BFV theory, possibly unify:

- ( -1 )-shifted 2-form attached to $M$ ( $B V$-form),
- 0-shifted 2-form attached to $\partial M$ (BFV-form).


## Outlook: the puzzle of quantisation

- Global geometric (non-perturbative)
- Focus on space of states



## Outlook: the puzzle of quantisation

- Setting to go beyond perturbative BV-BRST theory
- Usually one would consider $\Omega^{*}(X, \mathfrak{g})$ with $L_{\infty}$-structure and take shifted cotangent bundle $T^{*}[-1] \Omega^{*}(X, \mathfrak{g})$
- We can consider $\operatorname{Bun}_{G}^{\nabla}(X):=\left[X, \mathbf{B} G_{\text {conn }}\right]$ (or some concretification of this), and take derived critical locus $\mathbb{R} \operatorname{Crit}(S)(M)$ for a given $S: \operatorname{Bun}_{G}^{\nabla}(X) \rightarrow \mathbb{R}$
$\Longrightarrow$ Global geometric generalisation of BV-BRST theory
- Setting to go beyond BV-quantisation
- [Bunk, Sämann, Szabo], [Fiorenza, Sati, Schreiber]: higher geometric prequantisation of $n$-plectic structures and prequantum bundle $n$-gerbes
- [Safronov]: geometric quantisation of derived symplectic structures in derived algebraic geometry via algebraic bundle $k$-gerbes
$\Longrightarrow$ Beyond BV-quantisation by "higher derived" geometric (pre)quantisation?

Thank you for your attention!

